

On some properties of the ratio of matrix elements $\left| \frac{V_{ub}}{V_{cb}} \right|$ related to the Fermi momentum (p_f)

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Abstract. The most probable value of the Fermi momentum (p_f) has been suggested in the framework of the statistical model and its dependence on the ratio of the CKM matrix elements $|V_{ub}/V_{cb}|$ has been investigated. It has been suggested that the aforesaid ratio should be considerably enhanced compared to the existing estimates so as to reproduce reasonable values of the Fermi momentum, p_f .

1 Introduction

Recently much attention has been paid on the B-meson decay as it is a potential source of measuring the Cabibbo-Kobayashi-Masakawa (CKM) matrix elements. The rare hadronic B-meson decays described by the process $b \rightarrow u$ ($b \rightarrow u\ell\nu$) or charmless B-meson decay is particularly important for the study of the CKM matrix elements. Till now the only experimental method to determine $|V_{ub}|$ is through the end point of the lepton energy spectrum of $b \rightarrow u$ decay. The ACCMM model [1] provides an excellent method to determine the end point behaviour of the spectrum of a heavy quark into a massless quark and a lepton which is used to study the $|V_{ub}/V_{cb}|$ where the Fermi momentum (p_f) is a very important free parameter in determining $|V_{ub}/V_{cb}|$. Hwang et al. [2,3], Choi et al. [4] have also investigated the dependence of V_{ub} , V_{cb} on p_f in the framework of the ACCMM model.

In the present paper we have estimated p_f in the framework of the statistical model which enables us to estimate it directly by experimentally observed quantities like the decay constant (f_B) and mass (m_B) of the B-meson. We have also investigated the dependence of $|V_{ub}/V_{cb}|$ on the Fermi momentum in the context of the ACCMM model. With the input of $|V_{ub}/V_{cb}|$ predicted from different models, p_f has been determined. An analysis of the results suggests that the most plausible value of $|V_{ub}/V_{cb}|$ should be more than the recent estimates and that p_f is $\simeq 0.61$ GeV.

2 Estimate of p_f

In the low energy limit of the heavy meson annihilation, the Van-Royen-Weisskoff formula gives the relation between the decay constant f_B (for the B-meson) and the

ground state wavefunction at the origin $\psi_B(0)$ as,

$$f_B^2 = \frac{12}{m_B} |\psi_B(0)|^2 \quad (1)$$

In the statistical model the ground state wavefunction for meson is given by [5],

$$|\psi(r)|^2 = \frac{315}{64\pi r_B^{9/2}} (r_B - r)^{3/2} \theta(r_B - r) \quad (2a)$$

where θ is a step function.

$$\text{At } r = 0, \quad |\psi_B(0)|^2 = 315/64\pi r_B^3 \quad (2b)$$

where r_B represents the radius or size parameter of the B-meson.

It is pertinent to recall here that in the statistical model [5], the number density of quark (n_q or antiquarks $n_{\bar{q}}$) also represents the probability density of a meson, i.e. $n_q(r) = |\psi_B(r)|^2$ so that \bar{n}_q , the average density of quarks (similar for antiquarks also) becomes

$$\bar{n}_q = \frac{1}{2} |\psi_B(0)|^2 = 315/128\pi r_B^3 \quad (3)$$

Hence the average number density would be related to the corresponding (average) Fermi momentum [6] by,

$$p_f^3/3\pi^2 = \bar{n}_q$$

$$\text{or, } p_f^3 = \frac{1}{2} |\psi_B(0)|^2 3\pi^2 \quad (4)$$

Combining (1), (3) and (4) we get,

$$p_f^3 = 3\pi^2 f_B^2 m_B/24 \quad (5a)$$

Table 1. Present estimates of r_B and p_f corresponding to different potentials [8]

| Potentials [8] | r_B (GeV ⁻¹) | p_f (GeV) |
|--|----------------------------|-------------|
| 1. $-\frac{\alpha_s}{r} + kr$ | 2.41 | 1.18 |
| 2. $A + Br^\alpha$ | 3.7 | 0.77 |
| 3. $\lambda \frac{(r^\alpha - 1)}{\alpha} + c$ | 3.7 | 0.77 |
| 4. $-\frac{\alpha_c}{r} + k^1 r$ | 3.5 | 0.81 |
| 5. $c \ln\left(\frac{r}{r_B}\right)$ | 3.65 | 0.78 |

and

$$p_f r_B = 2 \cdot 85 \quad (5b)$$

with $f_B = (0.19 \pm 0.04)$ GeV [7] and $m_B = 5.28$ GeV [3], we get $p_f = 0.61$ GeV from (5a) which is in good agreement with the recent estimates of Hwang et al. ($p_f = 0.54^{+0.16}_{-0.15}$) GeV [3] and Choi et al. [4] ($p_f = 0.68$ GeV).

The expression (5b) represents a relation between Fermi momentum and the radius parameter of the B-meson and it measures the Fermi momentum of the two body bound state. In order to calculate p_f from (5b), the knowledge of r_B is essential. We have estimated r_B from the relativistic Hamiltonian of heavy meson by minimizing it. The values of p_f are thus obtained with the input of different r_B corresponding to different potentials [8]. The results are displayed in Table 1. It is observed that $r_B = 0.74$ fm (3.7 GeV) corresponding to $p_f = 0.77$ GeV agrees fairly well with the estimates of [3, 4].

3 Determination of p_f (ACCMM model)

The rare decay of B-meson transition via $b \rightarrow ul\nu$ contributes to the end point energy of lepton spectrum. It has been explicitly shown that the $b \rightarrow ul\nu$ transition is responsible for the excess of leptons with momenta above the kinematical limit for $b \rightarrow el\nu$ transition. In the ACCMM model the absolute value of $|V_{ub}|$ is determined from the behaviour of spectrum near the end point and $|V_{ub}/V_{cb}|$ is expressed as a function of the Fermi momentum p_f .

In the ACCMM [1] model the heavy quark (b) is treated as a virtual particle of invariant mass ‘ W ’ such that (with antiquark as the spectator),

$$W^2 = M_B^2 + m_{sp}^2 - 2m_B \sqrt{p^2 + m_{sp}^2} \quad (6)$$

where M_B is the mass of the B-meson. ‘ m_{sp} ’ represents the mass of the spectator antiquark and ‘ p ’ is the momentum of the ‘ b ’ quark inside B-meson.

Assuming the momentum distribution of the virtual ‘ b ’ quark inside the meson to be of Gaussian type we have,

$$\psi_B(p) = \frac{4}{\sqrt{\pi} p_f^3} e^{-p^2/p_f^2} \quad (7)$$

The lepton energy spectrum of the B-meson decay is given by,

$$d\Gamma_B/dE_l = \int_0^{p_f} dp p^2 \psi(p) \quad d\Gamma_b/dE_l \quad (8)$$

Table 2. Estimate of p_f for different values of $|V_{ub}/V_{cb}|$

| Ref. | $ V_{ub}/V_{cb} $ | p_f (GeV) |
|--------------|-------------------|-------------|
| ACCMM [1] | 0.1 | 0.297 |
| ARGUS [12] | 0.08 ± 0.02 | 0.326 |
| Melikov [11] | 0.108 ± 0.02 | 0.319 |
| Wirbel [10] | 0.3 | 0.870 |
| Present | 0.2 | 0.61 |

For the evaluation of $(d\Gamma_B/dE_l)$, we have parameterized $d\Gamma_b/dE_l$ by, (for large ‘ p ’),

$$d\Gamma_b/dE_l = A/p^2 + B \quad (9)$$

where $A = 0.285$, $B = 0.049$.

From (8) and (9) we get,

$$d\Gamma_B/dE_l = \frac{1.23}{p_f^2} - 0.085 \quad (10)$$

The decay width $\tilde{\Gamma}(p_f)$ has been estimated at the end point spectrum of the b-meson decay in the range $2.3 < E_l < 2.6$ where the contribution is due to the rare ‘ b ’ decay only so that defining $\tilde{\Gamma}(p_f)$ by,

$$\tilde{\Gamma}(p_f) = \int_{2.3}^{2.6} \frac{d\Gamma_B}{dE_l} \cdot E_l$$

we get,

$$\tilde{\Gamma}(p_f) = \frac{0.369}{p_f^2} - 0.025 \quad (11)$$

The experimentally measured width Γ_{expt} is given by,

$$\Gamma_{\text{expt}} = |V_{ub}|^2 \cdot \tilde{\Gamma}(p_f) \quad (12)$$

As Γ_{total} is proportional to $|V_{ub}/V_{cb}|^2$, hence we arrived at,

$$\Gamma_{\text{expt}}/\Gamma_{\text{total}} \propto \left| \frac{V_{ub}}{V_{cb}} \right|^2 \cdot \tilde{\Gamma}(p_f) \quad (13)$$

or,

$$\tilde{\Gamma}(p_f) = |V_{ub}/V_{cb}|^2 \cdot \tilde{\Gamma}(p_f)_{p_f=0.3} / \left| \frac{V_{ub}}{V_{cb}} \right|^2 \quad (14)$$

We have used $|V_{ub}/V_{cb}|^2$ from Isgur [9] and $\tilde{\Gamma}(p_f)$ at $p_f = 0.3$ GeV. With the input of $|V_{ub}/V_{cb}|$ suggested in different models we have estimated p_f using (14). The results are displayed in Table 2. It is found that the existing estimates of $|V_{ub}/V_{cb}|$ generate very low values of p_f except the value given by Wirbel et al. [10].

4 Results and discussions

In the present paper we have estimated p_f in the framework of the statistical model. The value of p_f obtained from (5a) with input of f_B and m_B agrees closely with the recent estimates [2–4]. As the decay constant (f_B) and m_B can be determined with considerable accuracy, the value of

p_f obtained using them (f_B, m_B) may not be far from the real value. Moreover (5b) gives a relation between Fermi momentum and radius of the two quark bound states. It is to be noted that the radius corresponding to a power law potential yields a reasonable value of p_f (0.77 GeV). Hwang et al. [2] have also observed that $p_f = 0.5$ GeV corresponds to the radius of the B-meson $r_B = 0.39$ fm, which is very low. They have used the Cornell potential to describe the interquark potential. In the present investigation we have found that $r_B = 0.74$ fm (corresponding to the power law type potential) yields reasonable estimate of p_f . So it may be suggested that the interacting potential between the constituent quark is fairly well described by the power law type potential.

Melikov [11] has analysed the semileptonic decays of the heavy meson within the dispersion formulation of the constituent quark model and have obtained $|V_{ub}/V_{cb}| = 0.108 \pm 0.02$. Although this estimate agrees with the recent value predicted by ARGUS ($|V_{ub}/V_{cb}| = 0.08 \pm 0.02$) [12], it produces low values of p_f like the estimate of [1]. On the other hand, Wirbel et al. [10] have assured that the ratio of CKM matrix elements has the bound $|V_{ub}/V_{cb}| \leq 0.3$ and it produces fairly large values of p_f towards its upper limit as is evident from Table 2. With our computed values of p_f ($= 0.61$ GeV) we get $|V_{ub}/V_{cb}| = 0.2$. So in the present investigation it has been observed that to generate reasonable values of p_f $|V_{ub}/V_{cb}|$ should have higher values than the existing estimates. We get $|V_{ub}/V_{cb}| = 0.2$ which is much more than the estimates by other workers [2–4]. It seems that the estimate predicted by Wirbel et al. [10] is more plausible (towards upper limit) and is supported by the present work. However it has been pointed out by ARGUS [12] that using only a restricted portion of the spectrum it is very difficult to extract the value of $|V_{ub}/V_{cb}|$. More experimental efforts are needed for a better estimate of $|V_{ub}/V_{cb}|$ which is a very important parameter for the understanding of CP-violation.

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